Chapter Seven - Vibrating Strings

The displacement u of a string is described by the equation

$$\frac{\partial}{\partial x}\left(T(x)\frac{\partial u}{\partial x}\right) - \rho(x)\frac{\partial^2 u}{\partial u^2} = 0,$$

where T(x) is the tension and $\rho(x)$ is the density. We have already seen this in the hanging chain problem—there the tension is proportional to x and the density is constant. Let's go back to the simpler problem of a uniform string fixed at the ends x = 0 and $x = \pi$. In this case the tension and the density are both constant: say T(x) = T and $\rho(x) = \rho$. Then

$$u_{xx} - \frac{\rho}{T}u_{tt} = 0, \ 0 < x < \pi$$

$$u(0,t) = u(\pi,t) = 0, \ \text{and}$$

$$u(x,0) = f(x), \ u_t(x,0) = g(x).$$

From our vast knowledge of eigenvalue problems we know to let $u = \sum_{n=1}^{n} \alpha_n(t) \sin nx$, which gives

$$\sum_{n=1}^{\infty} \left[-n^2 \alpha_n(t) - \frac{\rho}{T} \alpha_n''(t) \right] \sin x = 0.$$

Thus

$$\alpha_n''(t)+n^2\frac{T}{\rho}\alpha_n=0,$$

which has solutions

$$\alpha_n(t) = a_n \cos nvt + b_n \sin nvt$$
, where
 $v = \sqrt{\frac{T}{\rho}}$.

Hence,

$$u(x,t) = \sum_{n=1}^{\infty} (a_n \cos nvt + b_n \sin nvt) \sin nx.$$

From the initial conditions, we know

$$a_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin nx dx$$
, and $b_n = \frac{2}{\pi nv} \int_0^{\pi} g(x) \sin nx dx$.

Observe that the solution *u* is periodic in *t*: $u(x, t) = u(x, t + 2\pi/v)$

Example. Suppose v = 1, g(x) = 0, and

$$f(x) = \begin{cases} x/2 & 0 \le x \le \pi/2 \\ -(x-\pi)/2 & 1/2 < x \le 1 \end{cases}$$

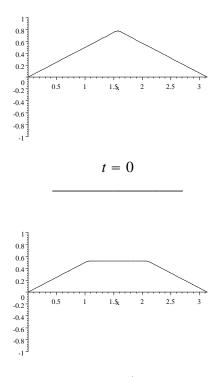
Then

$$a_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin nx dx = \frac{2 \sin \frac{n\pi}{2}}{\pi n^2},$$

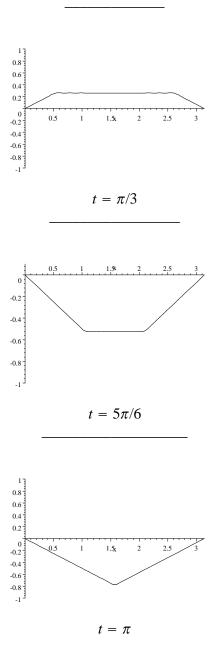
and $b_n = 0$. Hence,

$$u(x,t) = \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{\sin \frac{n\pi}{2}}{n^2} \cos nt \sin nx.$$

Let's see what this looks like for a sequence of values of time *t*.



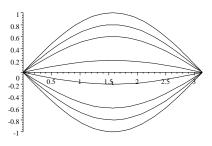
 $t=\pi/6$



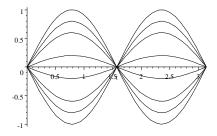
Suppose all the terms in the series for u save the one for n = 1 are zero. The solution then is simply

$$u(x,t) = (a_1 \cos vt + b_1 \sin vt) \sin x$$
$$= A \cos(vt + \varphi) \sin x.$$

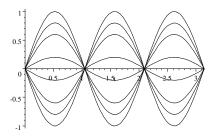
In this oscillation, the string always has the shape of a single arch of the sine curve and vibrates with the radian frequency v:



The solution for all n = 0 except n = 2 has the form $A\cos(2vt + \varphi)\sin 2x$. Here the string has the shape of $\sin 2x$ and oscillates with a frequency 2v, or twice the frequency of the previous solution:



Musically, this oscillation would sound an octave higher than the first one. Convince yourself that for n = 3, the string would vibrate with frequency 3v:



What is the musical interval between this one and the previous one?

The solutions like these in which all but one of values of n are zero are called **vibration modes**. Thus every solution is a superposition, or "sum", of these modes of vibration.

Let's consider again the solution in which g(x) = 0. Then

$$u(x,t) = \sum_{n=1}^{\infty} a_n \cos nvt \sin nx.$$

Drawing upon our vast knowledge of trigonometry, we see that

$$\cos nvt \sin nx = \frac{1}{2} [\sin(n(x+vt)) + \sin(n(x-vt))].$$

Hence,

$$u(x,y) = \frac{1}{2} \sum_{n=1}^{\infty} a_n [\sin(n(x+vt)) + \sin(n(x-vt))]$$

= $\frac{1}{2} \left(\sum_{n=1}^{\infty} a_n \sin(n(x+vt)) + \sum_{n=1}^{\infty} a_n \sin(n(x-vt)) \right).$

Now we know that for most all x

$$\widetilde{f}(x) = \sum_{n=1}^{\infty} a_n \sin nx,$$

where \tilde{f} is the odd periodic extension of f. Thus

$$\widetilde{f}(x+vt) = \sum_{n=1}^{\infty} a_n \sin(n(x+vt)), \text{ and } \widetilde{f}(x-vt) = \sum_{n=1}^{\infty} a_n \sin(n(x-vt)),$$

and our solution is

$$u(x,t) = \frac{1}{2} \left(\widetilde{f}(x+vt) + \widetilde{f}(x-vt) \right)$$

This is special case of what is called **D'Alembert's** formula of the wave equation. In case g is not zero, there is the full version of D'Alembert's solution:

$$u(x,t) = \frac{1}{2} \left(\widetilde{f}(x+vt) + \widetilde{f}(x-vt) \right) + \frac{1}{2v} \int_{x-vt}^{x+vt} g(s) ds$$